

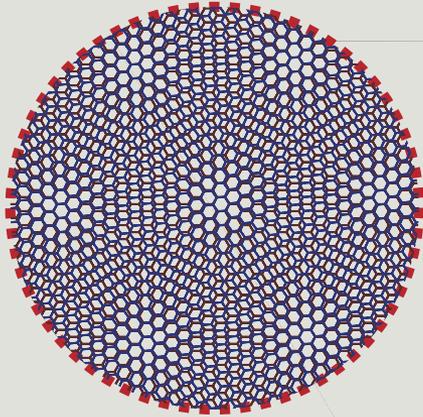
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## Quantum Anomalous Hall Effect in twisted bilayer graphene

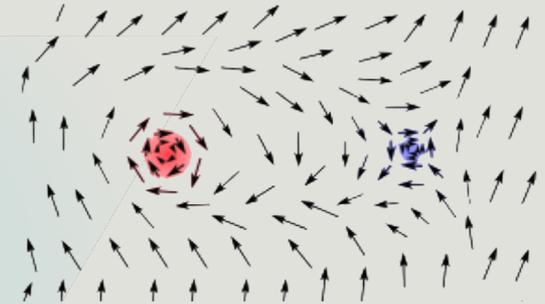
FLAT CLUB, LOUK RADEMAKER, 11 OCTOBER 2019

# Context & Relevance

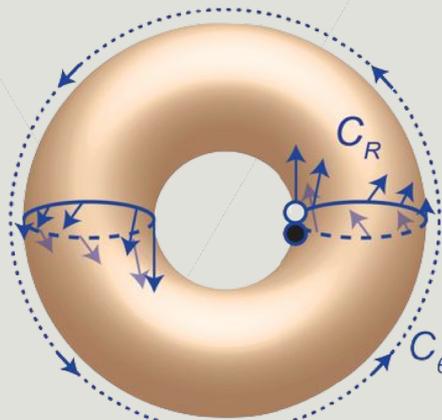
Moiré patterns  
in Twisted Heterostructures



Two-dimensional  
Magnetism

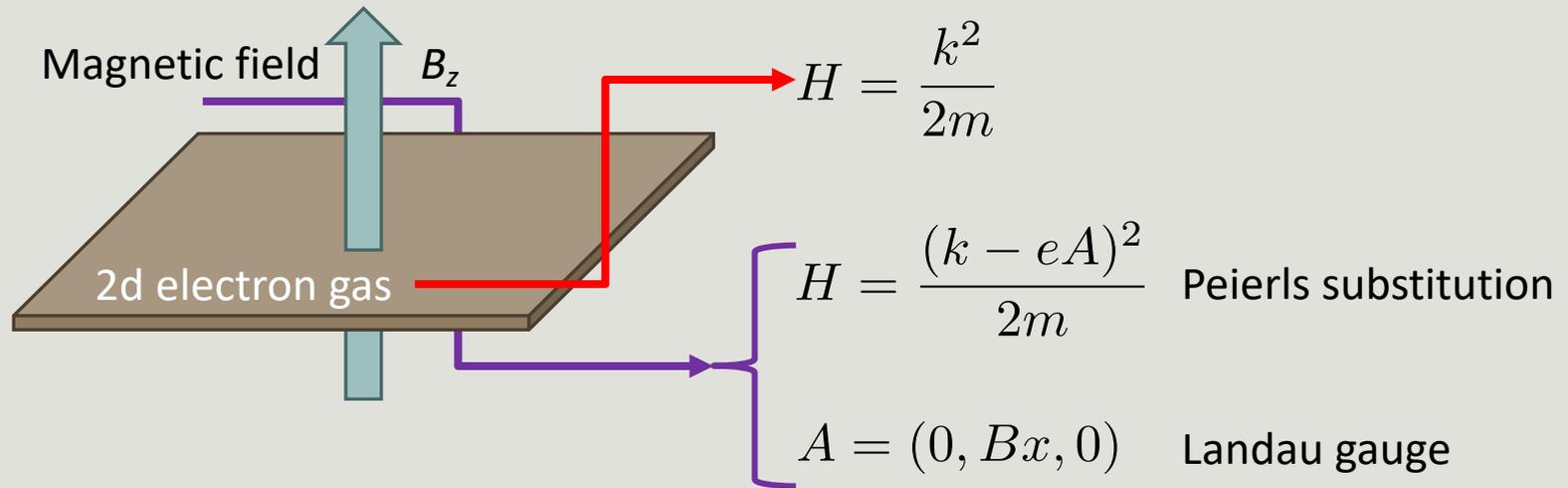


Quantum Anomalous  
Hall effect in  
Twisted Bilayer graphene



Band  
Topology

# Quantum Hall effect



$$H = \frac{\mathbf{k}_x^2}{2m} + \frac{1}{2m} (k_y - eBx)^2$$

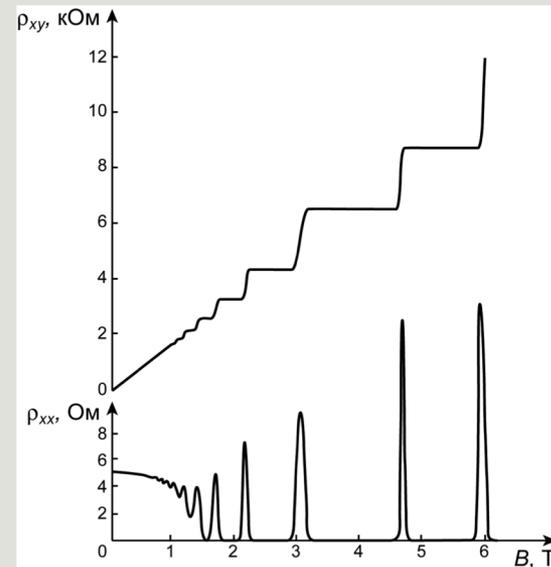
Shifted harmonic oscillator for each  $k_y$

$$H = \omega_C \left( \mathbf{n} + \frac{1}{2} \right)$$

Landau levels

$$\omega_C = \frac{eB}{m}$$

Cyclotron frequency

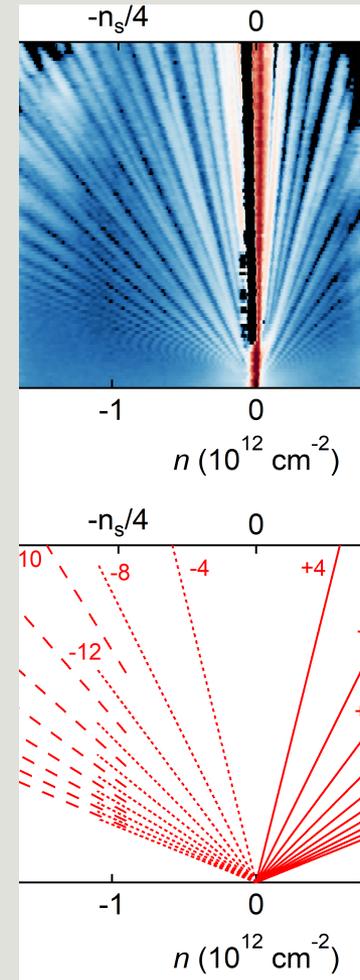
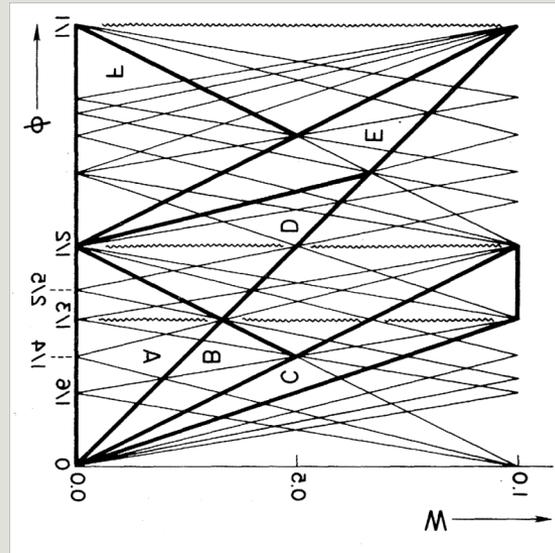
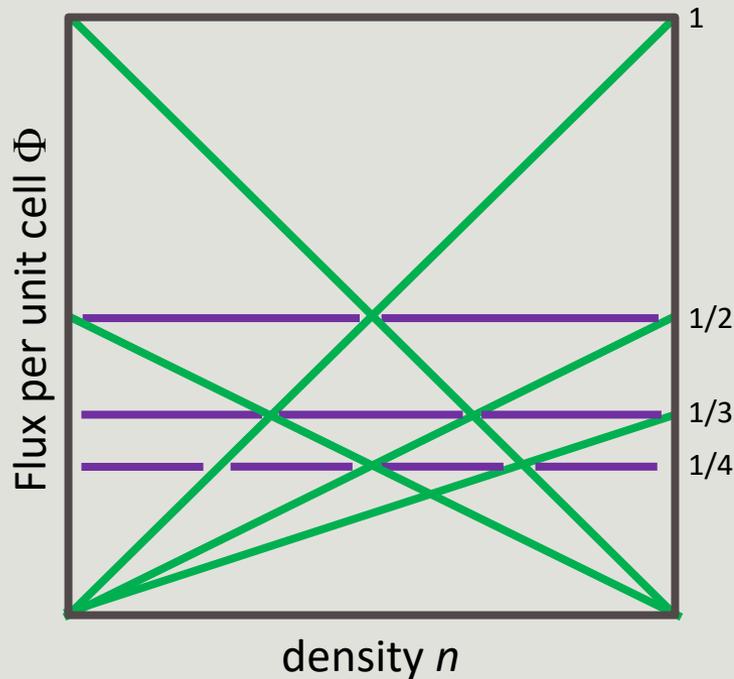


# Field vs. density

Number of states per Landau level:  $\Phi/\Phi_0$

↳ Insulating state when filled a Landau level

Claro-Wannier diagram (PRB 1979)



One  $\Phi_0$  per unit cell: { 40,000 T for graphene  
16 T for twisted bilayer graphene

Yankowitz  
Science 2019

# Streda formula

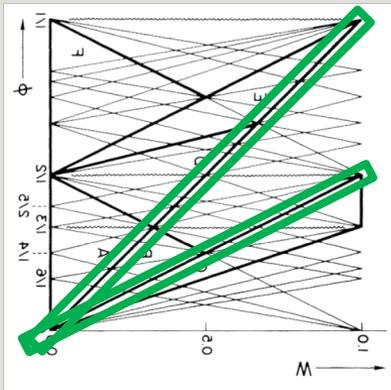
Hall response of an insulator:  $\sigma = \begin{pmatrix} 0 & \sigma_{xy} \\ \sigma_{xy} & 0 \end{pmatrix}$  with  $j = \sigma E$

Use the **continuity** & **Maxwell-Faraday** equations:  
*(practically: tune B field and gate voltage such that you remain insulating)*

$$\frac{\partial n}{\partial t} = -\nabla \cdot j = -\sigma_{xy}(\partial_x E_y - \partial_y E_x) = \sigma_{xy} \frac{\partial B}{\partial t}$$

Gives us the **Streda formula**

$$\sigma_{xy} = \frac{\partial n}{\partial B}$$



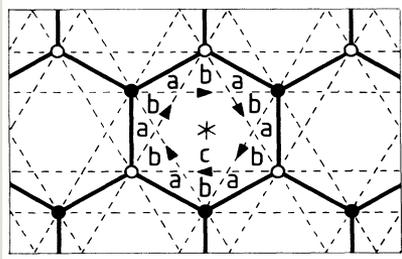
$$\Phi / \Phi_0 = \frac{1}{\nu} n / A$$

$$\sigma_{xy} = \nu \frac{e^2}{h}$$

Integer Quantum Hall Effect

# Haldane model (1988)

Can you have **QHE** without a magnetic field on a **lattice**?



$$t_{ij}c_i^\dagger c_j \rightarrow t_{ij}c_i^\dagger c_j e^{i \int_j^i \vec{A} \cdot d\vec{r}}$$

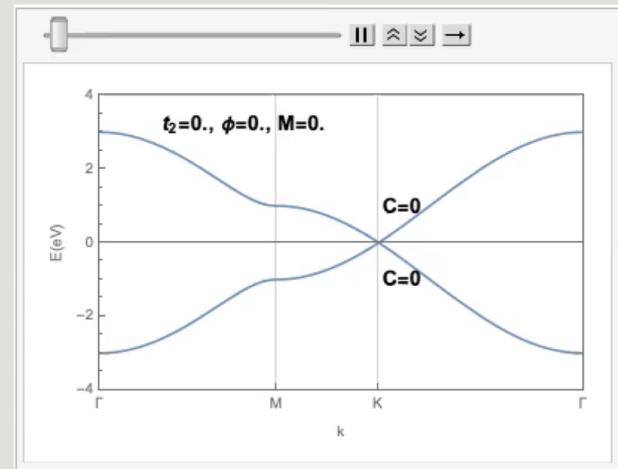
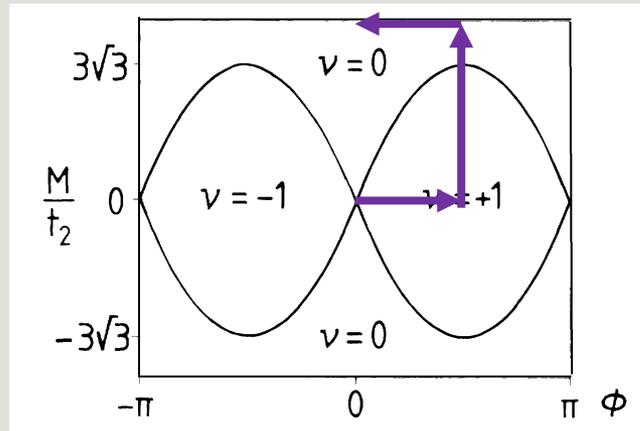
**No net flux** but the next-nearest-neighbor hopping acquires a phase

$$\mathbf{H}(\mathbf{k}) = 2t_2 \cos\phi \left( \sum_i \cos(\mathbf{k} \cdot \mathbf{b}_i) \right) \mathbf{I} + t_1 \left[ \sum_i [\cos(\mathbf{k} \cdot \mathbf{a}_i) \sigma^1 + \sin(\mathbf{k} \cdot \mathbf{a}_i) \sigma^2] \right] + \left[ M - 2t_2 \sin\phi \left( \sum_i \sin(\mathbf{k} \cdot \mathbf{b}_i) \right) \right] \sigma^3$$

NNN hopping  
with phase

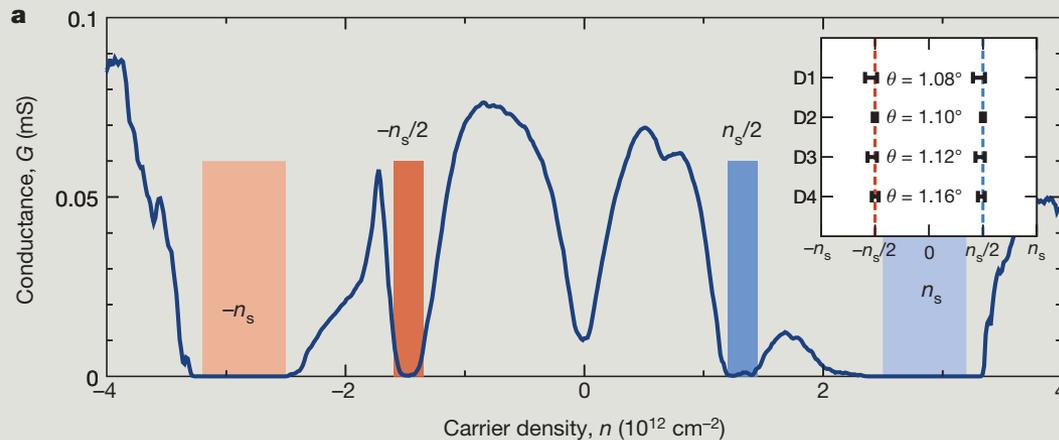
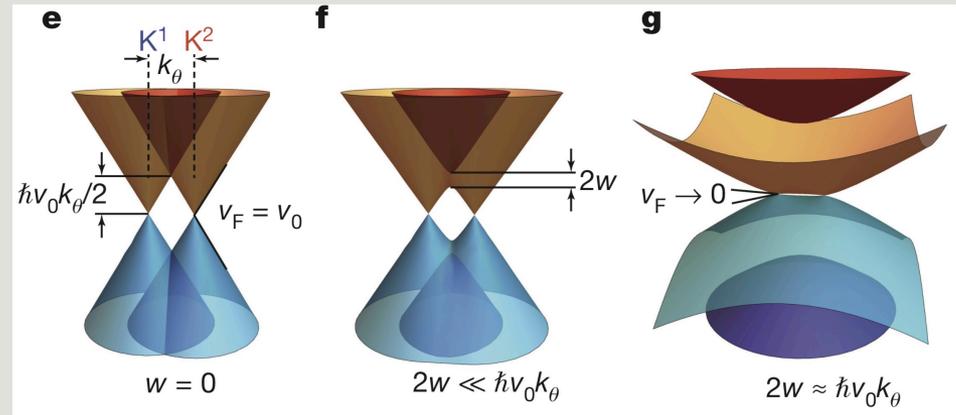
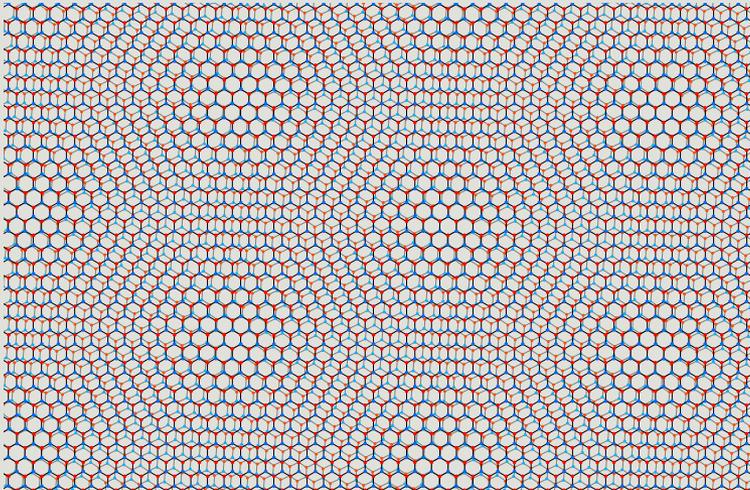
“graphene”

trivial gap



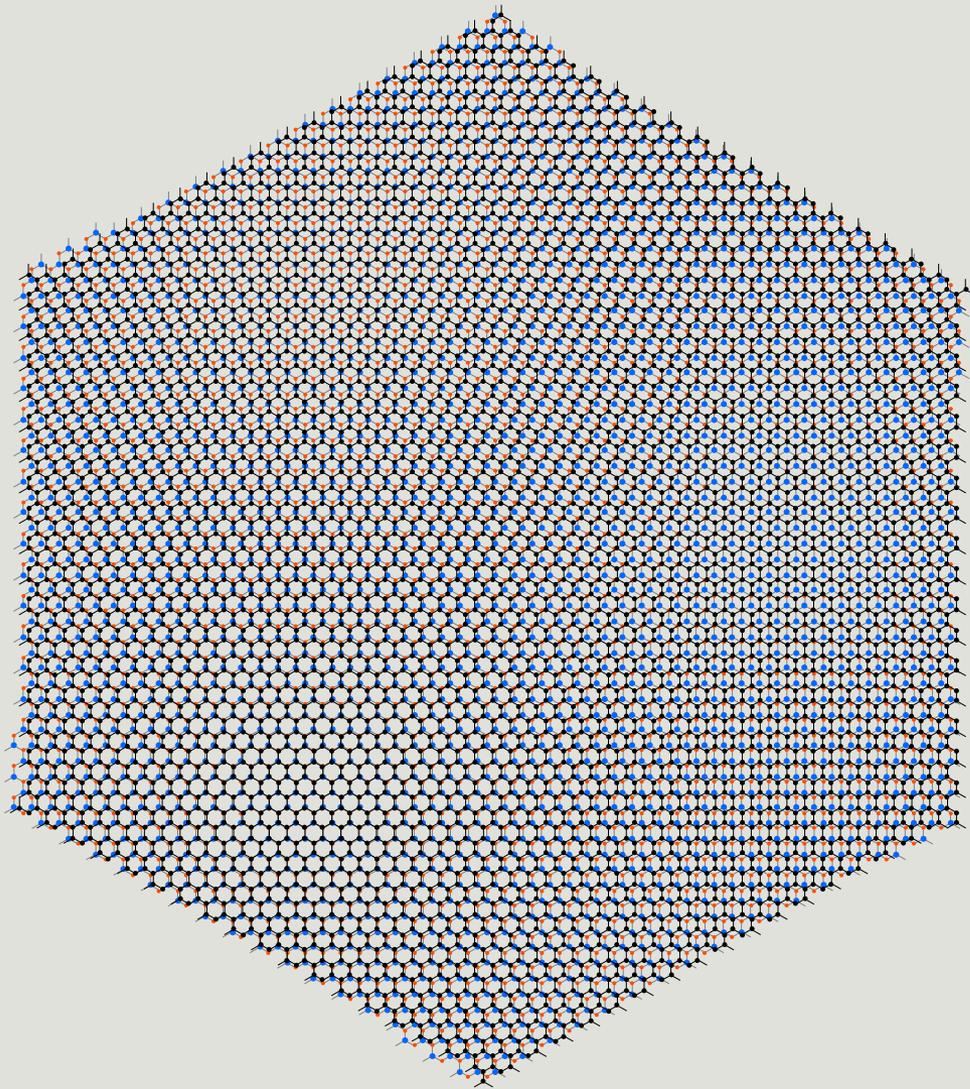
# Twisted Bilayers

Quick recap of twisted bilayer physics



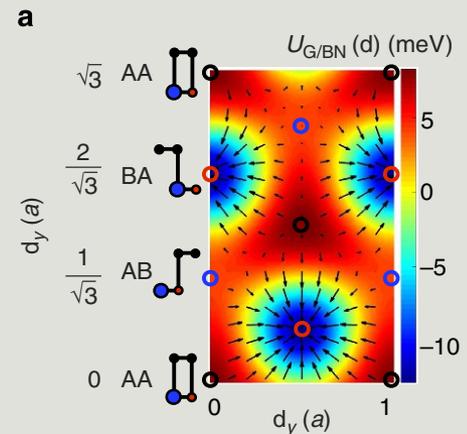
(Cao Nature 2018)

# hBN alignment with graphene



hBN substrate is hexagonal but **has 1.6% larger lattice constant**

When **graphene** is placed on top, structural relaxation favors BA stacking



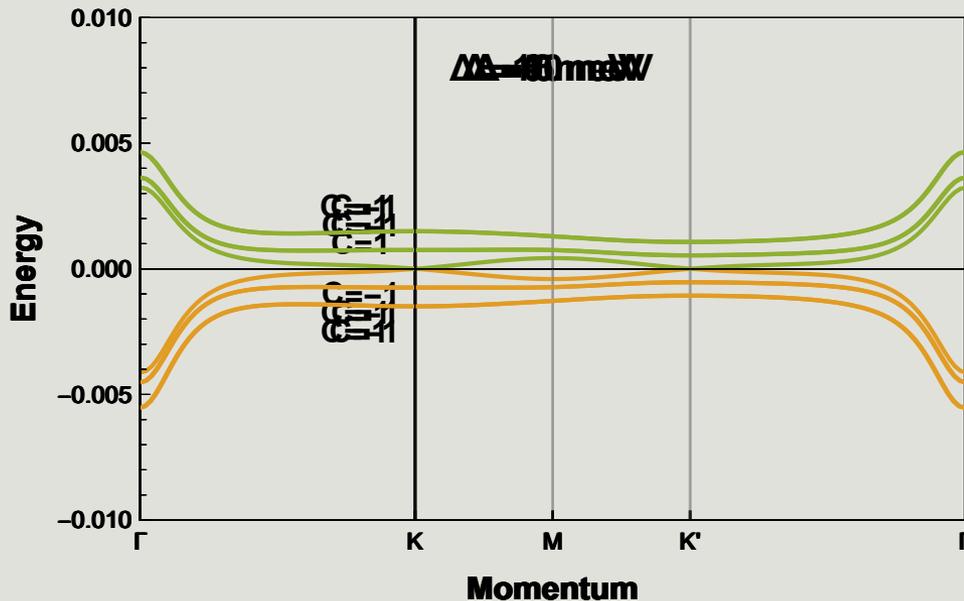
Leading in graphene to **sublattice symmetry breaking**

(Jung Nature Comm. 2015)

# hBN alignment on TBG

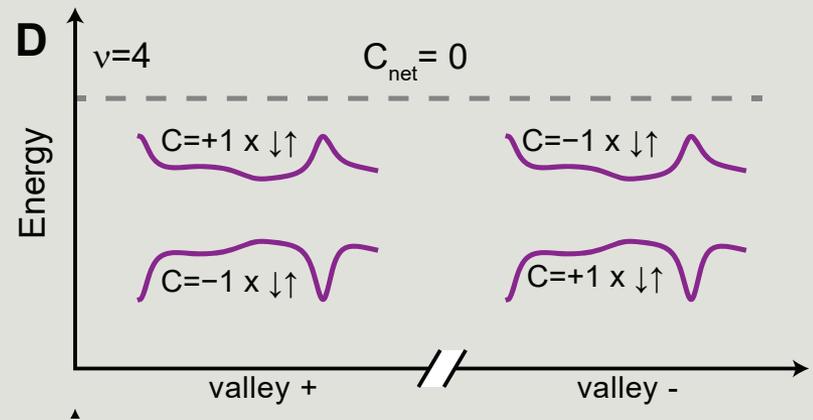
Very beautiful result for TBG:

- Include the sublattice breaking **on one of the layers**
- Leads to a **gap** at K and K' point



Above: bandstructure for a single valley.

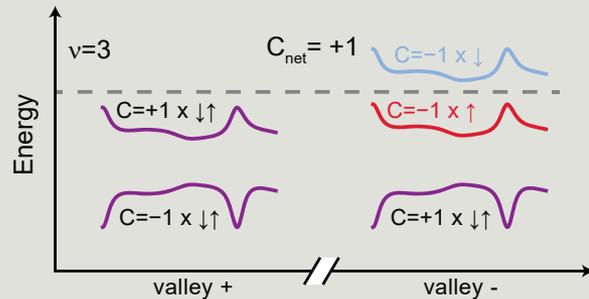
(Bultinck arXiv:1901.08110)



(Serlin arXiv:1907.00261)

# Ferromagnetism

Imagine at  $\nu=3$  filling you have a **spin & valley polarized state**



Insulator

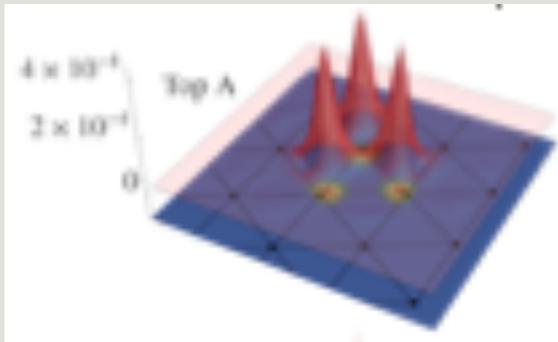
Net nonzero Chern number

Breaks time-reversal symmetry

**Quantum Anomalous Hall effect!**

But **why would you expect ferromagnetism?**

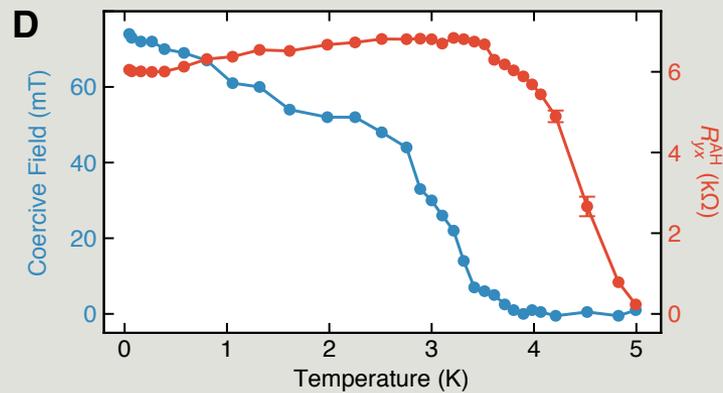
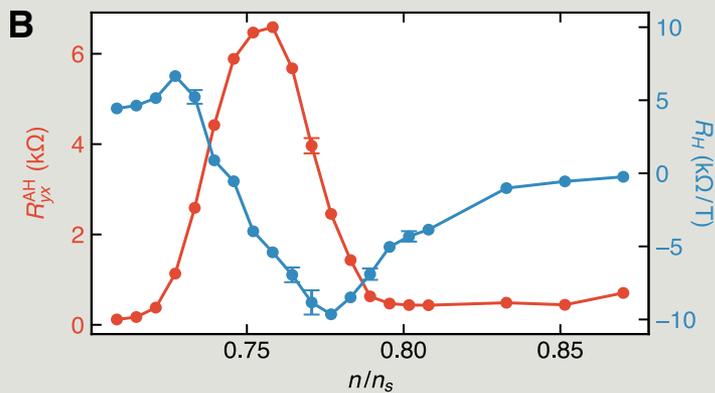
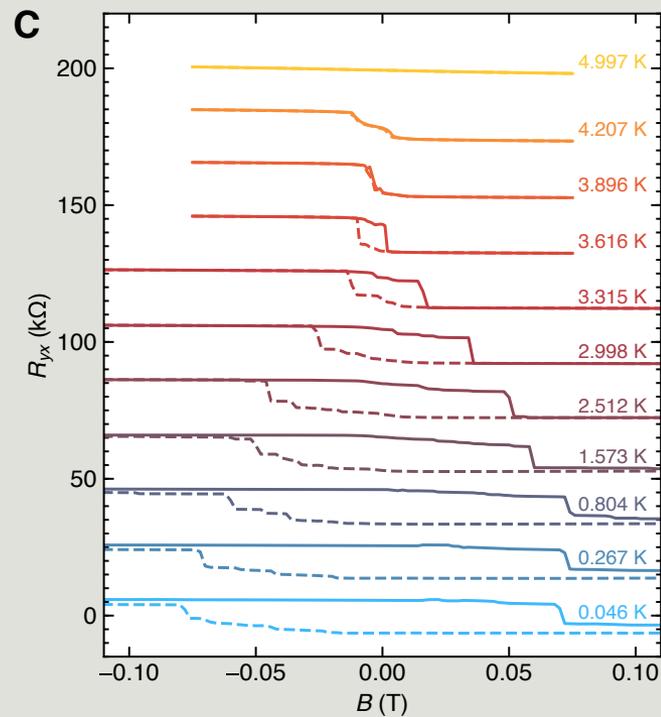
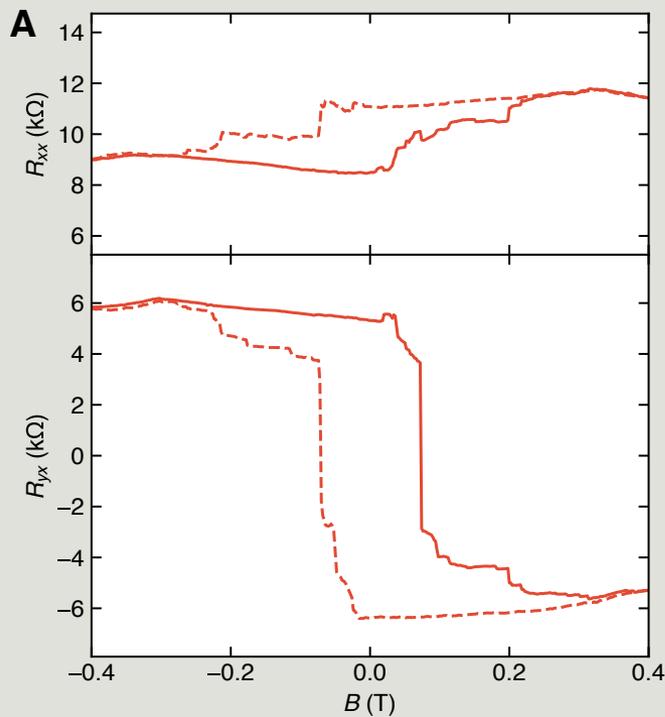
- Superexchange typically gives antiferromagnetic coupling
- **Direct exchange** between **Wannier orbitals** gives ferromagnetic coupling



(Ochi, arXiv:1805.09606; Kang, arXiv:1810.08642)

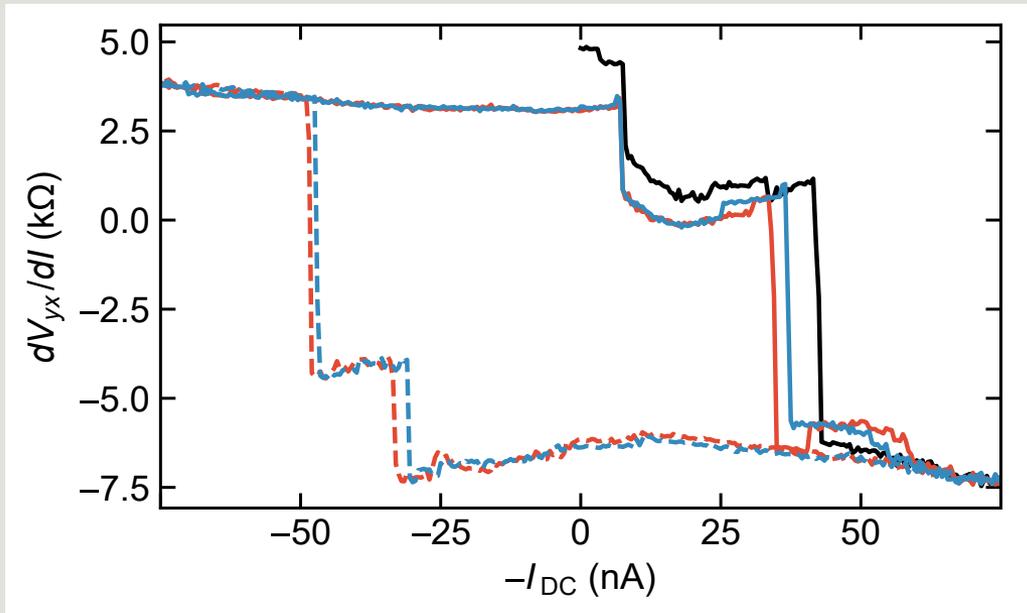
# Stanford Experiment (1)

(Sharpe arXiv:1901.03520)



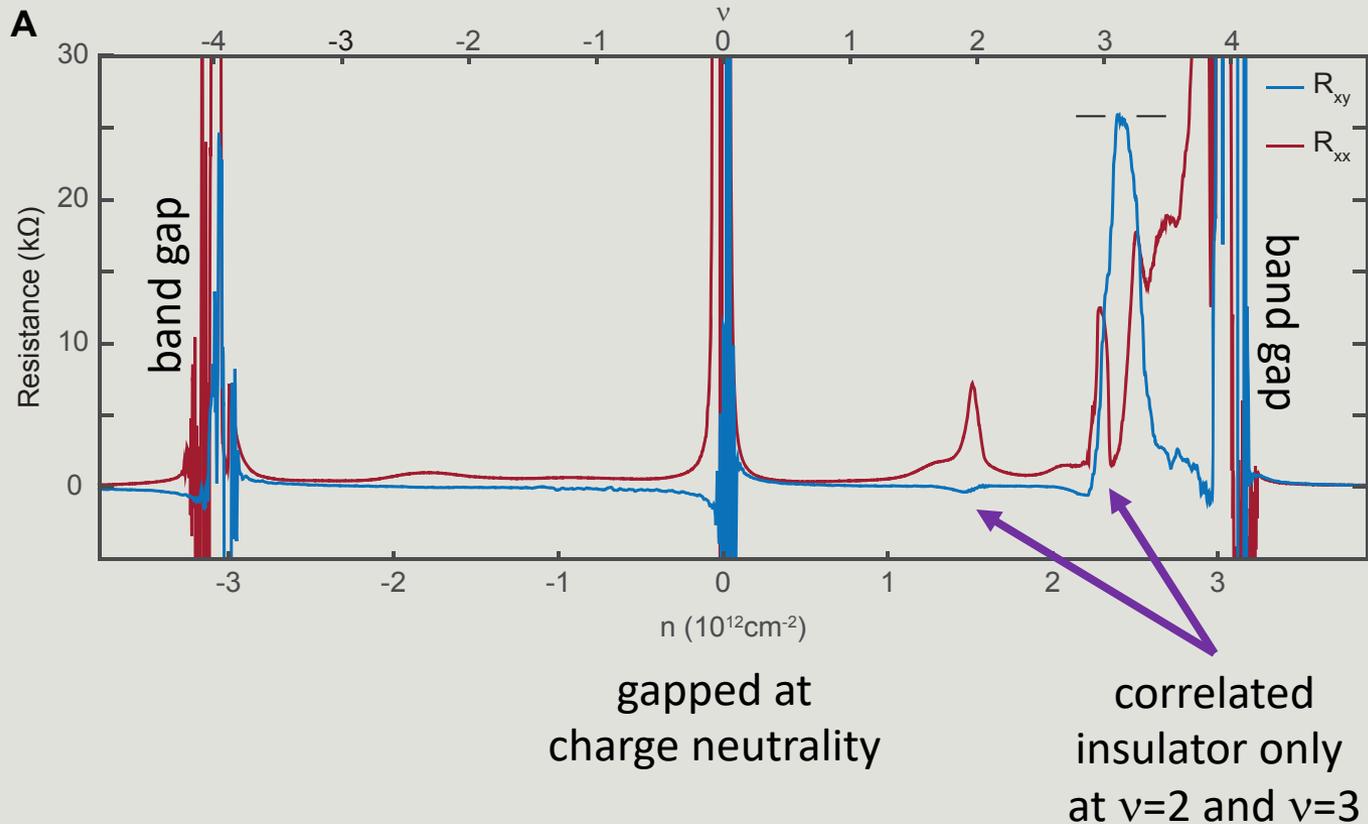
# Stanford Experiment (2)

Applying a **DC current** in the plane can switch the polarized state



# Santa Barbara Experiment (1)

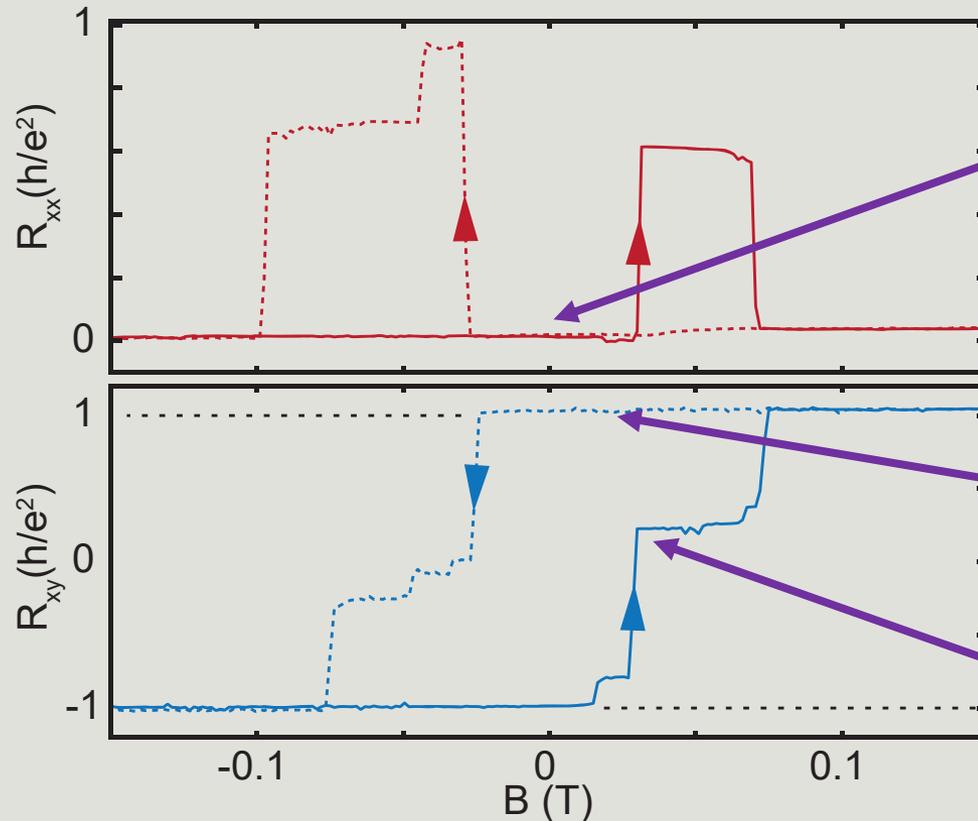
Measure **longitudinal and transverse** resistivity as a function of **doping**  
(done at weak field of 150 mT and low  $T=1.6$  K)



# Santa Barbara Experiment (2)

At  $\nu=3$  filling and  $T=1.6$  K, get **hysteresis** and a perfect **Quantum Anomalous Hall effect**

**B**



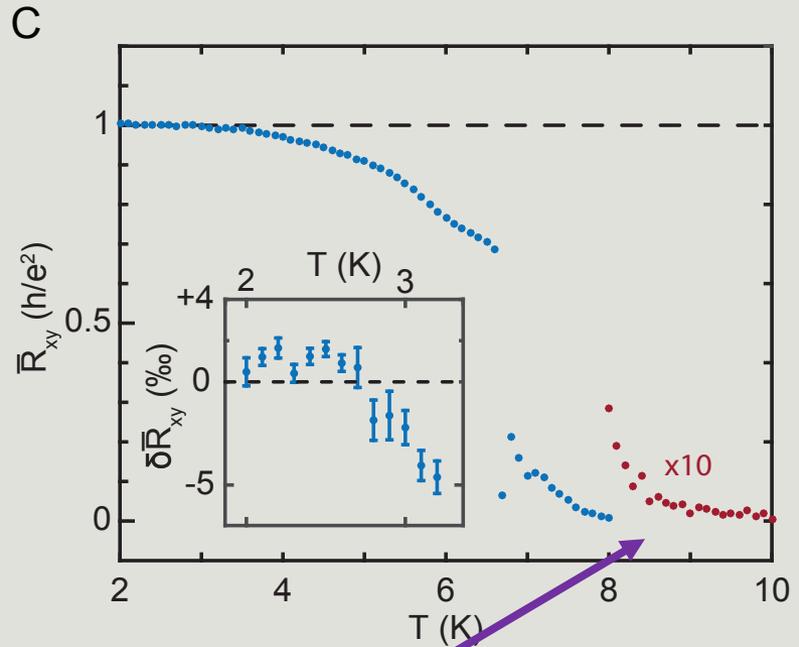
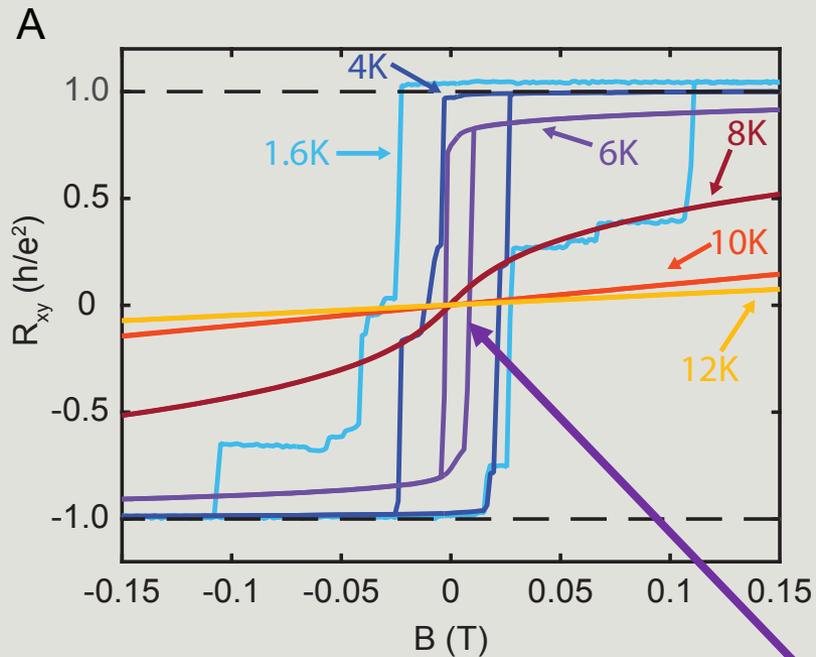
Really insulating at zero field

Quantized to within 0.1% of  $h/e^2$

Steps are likely caused by ferromagnetic domains

# Santa Barbara Experiment (3)

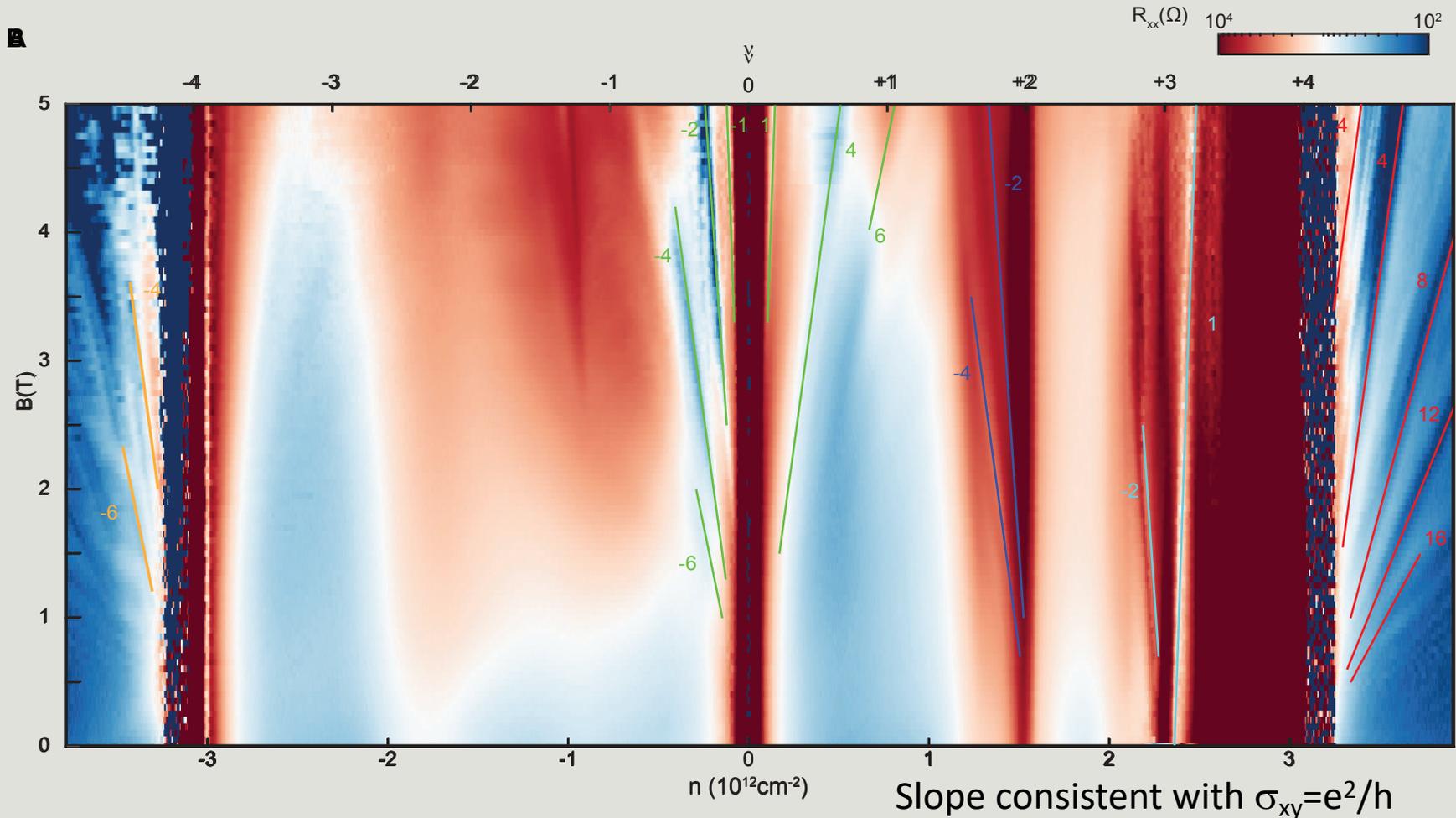
Same data, now as a function of **temperature**



Hysteresis loop disappears around 8-9 Kelvin

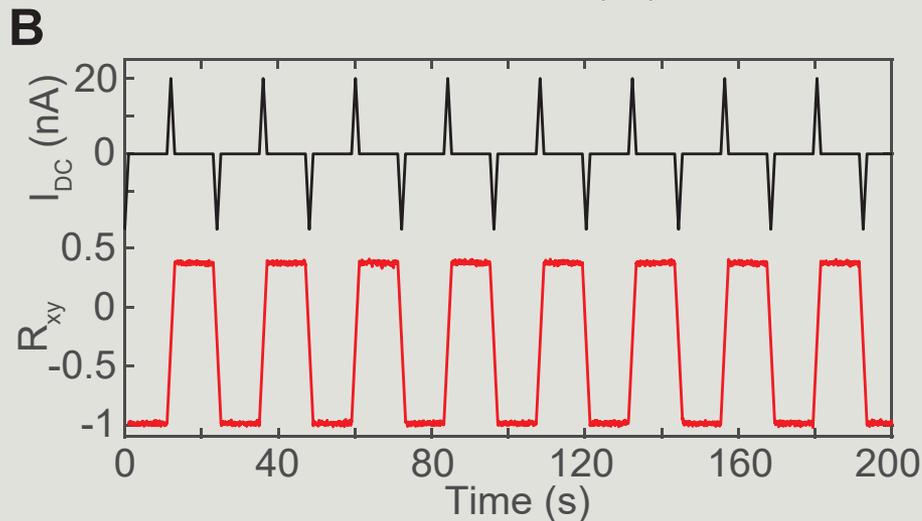
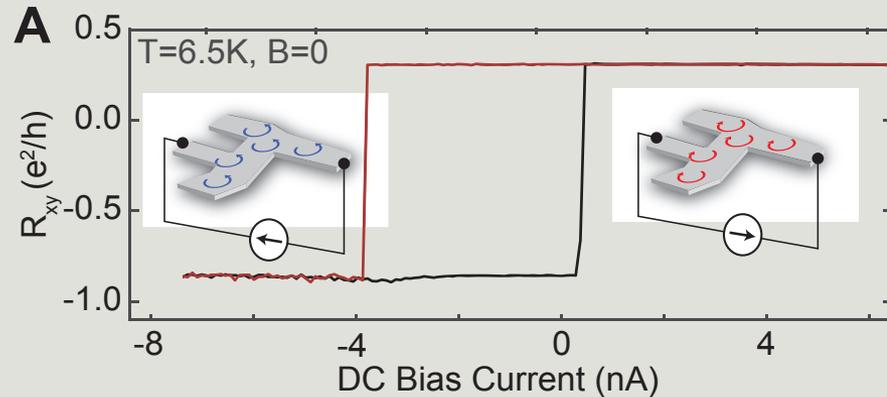
# Santa Barbara Experiment (4)

Same data, now as a function of **doping and field** to see Landau fans



# Santa Barbara Experiment (5)

Current-induced **switching** of magnetization at higher T



Note the offset in  $R_{xy}$ , because measurements are not “Onsager symmetrized” (?)

# Open questions

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- Why is the state spin/valley ferromagnetic?
- Why only the  $\nu=3$  filling – what about others?
- Why the current-switching of magnetization?
- Why does this effect only happen with aligned hBN?
- Why is there only QAH in California?